# Calculation of the Molecular Polarizability Tensor 

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#### Abstract

A semiempirical method is presented to compute the components of molecular polarizability with atomic polarizability components and the dipole tensor. Anisotropic atomic polarizabilities are oriented along the bonds and in the direction of all nonbonded atoms with a distance-dependent function. The average of the trace of each atomic polarizability matrix is the empirical result, $\alpha_{\mathrm{A}}(\mathrm{hhp})$, for atom A calculated by K. J. Miller. ${ }^{10}$ The dipole tensor is modified by omitting the terms connecting bonded atoms and damping the remaining ones to include the dipole approximation with increasing distance. The anisotropic atomic polarizabilites incorporate the nearest neighbor effects along the bonds and replace the dipole terms that overestimate contributions to the polarizability. The components of the molecular polarizability are scaled so that the average of the trace is equal to the empirical molecular polarizability, $\alpha(\mathrm{ahp})$. The results are invariant to rotations or translations of the coordinate system. A set of parameters is presented that yields a comparison between the scaled components of molecular polarizability and experimental results for approximately 120 molecules with a standard deviation of $3.5 \%$. The method is simple to use and suitable for molecular mechanics calculations.


## Introduction

Silberstein ${ }^{1-3}$ proposed the theory of the atom dipole approximation. It was further developed by Rowell and Stein, ${ }^{4}$ Mortensen, ${ }^{5}$ DeVoe, ${ }^{6}$ Applequist, ${ }^{7}$ Birge, ${ }^{8}$ and Thole. ${ }^{9}$ If the atomic polarizability tensor for each atom $i$ is given by $\alpha_{i}$, a $3 \times 3$ matrix, then the induced dipole moment $\mu_{i}$ at each atom $i=1,2,3, \ldots$ $L$ is given by

$$
\begin{equation*}
\mu_{i}=\alpha_{i}\left[\mathbf{E}_{i}-\sum_{j(\neq i)} \mathbf{T}_{i j} \mu_{j}\right] \quad i, j=1,2,3, \ldots, L \tag{1}
\end{equation*}
$$

where $\mu_{i}$ and $\mathbf{E}_{i}$ (the applied electric field) are column vectors containing three elements, the $x, y$, and $z$ components for atom $i$, and $\mathbf{T}_{i j}$ ( $3 \times 3$ matrices) are the dipole tensors connecting atoms $i$ to $j$

$$
\begin{align*}
& T_{p i q j}=-\left[\partial^{2} / \partial p_{i} \partial q_{j}\right]\left[1 / r_{i j}\right]  \tag{2}\\
& =\left[-1 / r_{i j}{ }^{5}\right]\left[3 p_{i} q_{j}-\delta_{p q} r_{i j}{ }^{2}\right]
\end{align*}
$$

where $p, q=x, y$, and $z$ are the Cartesian components. The $\alpha_{i}$ and $\mathrm{T}_{i j}$ matrices will be modified in this paper in a theoretical development of a modified dipole approximation. The term in brackets of eq 1 is the total electric field. It is a sum of the applied plus the induced electric fields. Equation I may be inverted to yield

$$
\begin{equation*}
\alpha_{i}^{-1} \mu_{i}+\sum_{j(\neq i)} \mathbf{T}_{i j} \mu_{j}=\mathbf{E}_{i} \tag{3}
\end{equation*}
$$

or

$$
\begin{equation*}
\mathbf{A}_{\mu}=\mathbf{E} \tag{4}
\end{equation*}
$$

where $\mathbf{A}$ is the $3 L \times 3 L$ matrix representation of $\alpha_{i}^{-1}$ and $\mathrm{T}_{i j}$, and $\mu$ and $\mathbf{E}$ are the $3 L \times 1$ column vectors of $\mu_{i}$ and $\mathbf{E}_{i}$ in eq 3, which can be expressed in the form

$$
\mathrm{A}_{\mu}=\left[\begin{array}{llll}
\alpha_{1}^{-1} & \mathrm{~T}_{12} & \ldots & \mathrm{~T}_{1 L}  \tag{5}\\
\mathrm{~T}_{21} & \alpha_{2}^{-1} & \ldots & \mathrm{~T}_{2 L} \\
\cdot & \cdot & & \cdot \\
\cdot & \cdot & \cdot & \cdot \\
\cdot & \cdot & \cdot \\
\mathrm{~T}_{L 1} & \mathrm{~T}_{L 2} & \ldots & \alpha_{L}^{-1}
\end{array}\right]\left[\begin{array}{l}
\mu_{1} \\
\mu_{2} \\
\cdot \\
\cdot \\
\mu_{L}
\end{array}\right]
$$

(1) Silberstein, L. S. L. Philos. Mag. 1917, 33, 92.
(2) Silberstein, L. S. L. Philos. Mag. 1917, 33, 215
(3) Silberstein, L. Philos. Mag. 1917, 33, 521.
(4) Rowell, R. L.; Stein, R. S. J. Chem. Phys. 1967, 47, 2985
(5) Mortensen, E. M. J. Chem. Phys. 1968, 49, 3732.
(6) Devoe, H. J. Chem. Phys. 1965, 43, 3199.
which is a matrix of $3 \times 3$ submatrices and $3 \times 1$ vectors for each atom $i$. The inverse $\mathrm{B}=\mathrm{A}^{-1}$ yields the induced dipole in the molecule from eq 4 as

$$
\begin{equation*}
\mu=\mathrm{BE} \tag{6}
\end{equation*}
$$

which contains $L$ equations for the induced dipole at each atom

$$
\begin{equation*}
\mu_{i}=\sum_{j} \mathbf{B}_{i j} \mathbf{E}_{j} \quad \mathbf{i}=1,2,3, \ldots, L \tag{7}
\end{equation*}
$$

If a constant external field is applied then $\mathbf{E}_{1}=\mathbf{E}_{2}=\ldots \mathbf{E}$. The total induced moment is a sum over $\mu_{i}$, and the molecular polarizability tensor becomes a sum over all $3 \times 3$ matrices $\mathrm{B}_{i j}$ which, after diagonalization, yields $\alpha_{\text {mol }}$ with principal components, $\alpha_{1}(\mathrm{~mol}), \alpha_{2}(\mathrm{~mol})$ and $\alpha_{3}(\mathrm{~mol})$, namely

$$
\begin{equation*}
\alpha_{\text {mol }}=\mathbf{D}\left(\sum_{i j} \mathbf{B}_{i j}\right) \mathbf{D}^{1} \tag{8}
\end{equation*}
$$

$$
=\left[\begin{array}{lll}
\alpha_{1}(\mathrm{~mol}) & \cdot & \cdot \\
\cdot & \alpha_{2}(\mathrm{~mol}) & \cdot \\
\cdot & \cdot & \alpha_{3}(\mathrm{~mol})
\end{array}\right]
$$

The eigenvectors, $\mathbf{D}$, define the principal axes of molecular polarizability, where $t$ is the transpose. The average molecular polarizability is the average of the trace $\operatorname{Tr}\left(\alpha_{\text {mol }}\right) / 3$ or

$$
\begin{equation*}
\alpha(\mathrm{mol})=\left[\alpha_{1}(\mathrm{~mol})+\alpha_{2}(\mathrm{~mol})+\alpha_{3}(\mathrm{~mol})\right] / 3 \tag{9}
\end{equation*}
$$

In the empirical methods to date, the set of atomic polarizabilities is adjusted until $\alpha(\mathrm{mol}), \alpha_{1}(\mathrm{~mol}), \alpha_{2}(\mathrm{~mol})$, and $\alpha_{3}(\mathrm{~mol})$ approximate the experimental values. In the present approach, average atomic polarizabilities are taken as the empirical set $\alpha_{A}(a h p)$, atomic anisotropies are introduced and adjusted, the dipole approximation is modified and adjusted, and the molecular polarizability matrix is scaled uniformly so that $\operatorname{Tr}\left(\alpha_{\text {mol }}\right) / 3 \rightarrow \alpha(\mathrm{ahp})$ and the normalized components approximate the experimental anisotropy. ${ }^{10}$ Within the point dipole approximation eq 1 is exact. The extent to which the dipole approximation is applicable to molecular systems may be understood by introducing successive approximations. First the approximations will be presented, and then the importance of each will be discussed by reporting effects of each approximation.

The determination of atomic polarizabilities by Applequist et al. ${ }^{7}$ that reproduced experimental molecular polarizabilities resulted from a set of abnormally small atomic polarizabilities when

[^0]compared to $\alpha_{A}$ (ahp). Introduction of atomic anisotropic effects by $\mathrm{DeVoe}^{6}$ and $\mathrm{Birge}^{8}$ improved the results. The use of a potential interaction to define and orient $\alpha_{i}$ by Birge ${ }^{8}$ and in greater detail by Thole ${ }^{9}$ improved the atomic and molecular polarizabilities to yield $2-7 \%$ agreement between theory and experiment. The fact that the dipole approximation is a long-range one is an important aspect of the present method.

## Isotropic Atomic Hybrid and Attenuated Dipole Tensor Approximations

In the present approach, three approximations were developed and used with eqs 3-9.
(1) The average atomic polarizability is constrained to the empirical values $\alpha_{A}(a h p)$ for atoms classified as atomic hybrids. ${ }^{10}$ Anisotropy is introduced to reproduce experimental results on diatomic and then polyatomic molecules and also to replace the dipole terms between bonded atoms. The atomic polarizability is given by
$\alpha_{i}(\mathrm{ahp})=\alpha_{i}{ }^{0}(\mathrm{ahp})+\sum_{j} l_{i j}{ }^{-1} s_{i j} \exp \left(-g_{i j}\right) \mathbf{V}_{i j} \Delta \alpha_{i}(\mathrm{ahp}) \mathbf{V}_{i j}{ }^{j}$
where $\alpha_{i}{ }^{\circ}(\mathrm{ahp})$ is the isotropic polarizability of atoms $i=1,2$, ..., $L, t$ is the transpose, and

$$
\begin{equation*}
\Delta \alpha_{i}(\mathrm{ahp})=\alpha_{i}(\mathrm{ahp})\left[3 \mathrm{f}_{i}-\mathrm{I}\right] \tag{II}
\end{equation*}
$$

where

$$
\mathbf{f}=\left[\begin{array}{ccc}
f_{i \|} & 0 & 0  \tag{12}\\
0 & f_{i \perp} & 0 \\
0 & 0 & f_{i \perp}
\end{array}\right]
$$

is the scaled anisotropic correction $\left(f_{i \|}+2 f_{i \perp}=1\right)$ where $f_{i| |}$ is the fractional contribution of the atomic polarizability along a bond or direction $i$ to $j ; f_{i \perp}$ is the fractional component perpendicular to the direction of a bond or $i$ to $j ; l_{i j}=l_{\mathrm{b}}$ is the number of atoms bonded to $i$, and $l_{i j}=l_{\mathrm{n}}$, the total number of atoms -$l_{\mathrm{b}}=L-l_{\mathrm{b}}$, is 1 plus the number of atoms not bonded to atom $i ; \mathbf{V}_{i j}$ is the unitary transformation, a $3 \times 3$ matrix, which transforms $\Delta \alpha_{i}$ from the local coordinate system $i$ to $j$ to the space-fixed system; $g_{i j}=0$ if $i$ is bonded to $j$ and $g_{i j}=c\left[r_{i j}-\right.$ $\left.d \rho_{i j}\right] / \rho_{i j}$ if $i$ is not bonded to $j$, where $c=2.54$ and $d=0.25$ were found to be most effective; $\rho_{i j}=\rho_{i}+\rho_{j}$ is the sum of theoretical van der Waals radii of atoms $i$ and $j$; and $s_{i j}=\rho_{i j} / r_{i j}$.

The weighting factors $l_{\mathrm{b}}$ and $l_{\mathrm{n}}$ deserve comment. The sum over atoms bonded to atom $i$ includes exactly $l_{\mathrm{b}}$ atoms, whereas the sum over atoms not bonded to atom $i$ is less than $l_{\mathrm{n}}$. For example, in ethane, $l_{\mathrm{n}}=7$ for each of six hydrogen atoms, whereas $l_{\mathrm{n}}=$ 4 for each of two carbon atoms. In benzene, $l_{\mathrm{n}}=11$ for each of six hydrogen atoms and $l_{\mathrm{n}}=9$ for each of six carbons. These weight factors were found empirically.

The sum over all atoms bonded to atom $i$ includes $l_{\mathrm{b}}$ terms, each weighted by $l_{\mathrm{b}}{ }^{-1}$. The average polarizability of atom $i$ remains $\alpha_{i}($ ahp $)$, but the atomic polarizability matrix is redirected along a resultant of all bond directions. The sum over all atoms not bonded to atom $i$ includes a different number of terms than $l_{\mathrm{n}}$. These terms make no contribution to the average atomic polarizability because the trace of $\Delta \alpha_{i}$ and $\mathbf{V}_{i j} \Delta \alpha_{i} \mathbf{V}_{i j}{ }^{f}$ vanishes so that the rotation of the anisotropic contribution $\Delta \alpha_{i}$ along the direction $i$ to $j$ continues to modify $\alpha_{i}$ but leaves its overall trace and average invariant. In summary

$$
\begin{align*}
\operatorname{Tr}\left[\mathbf{V}_{i j} \Delta \alpha_{i} \mathrm{~V}_{i j}\right] & =\operatorname{Tr}\left[\Delta \alpha_{i}\right]=0  \tag{13}\\
\operatorname{Tr}\left(\alpha_{i}{ }^{\circ}\right) / 3 & =\alpha_{i}(\mathrm{ahp}) \tag{14}
\end{align*}
$$

and

$$
\begin{equation*}
\operatorname{Tr}\left(\alpha_{i}\right) / 3=\alpha_{i}(\mathrm{ahp}) \tag{15}
\end{equation*}
$$

follows from eqs 10-12. These changes in $\alpha_{i}$ occur mainly along the bonded atoms because those along the nonbonded directions $i$ to $j$ are damped by $\exp \left(-g_{i j}\right)$.

Table I. Atomic Hybrid Polarizabilities and Scaled Components ${ }^{a}$

| atomic <br> hybrid | $\alpha_{A}, \AA^{3}$ | $f_{1}, f_{1}$ | $f_{2}, f_{\perp}$ | $f_{3}, f_{\perp}$ |
| :--- | :---: | :---: | :---: | :---: |
| H | 0.387 | 0.3940 | 0.3030 | 0.3030 |
| F | 0.296 | 0.4600 | 0.2700 | 0.2700 |
| Cl | 2.315 | 0.4400 | 0.2800 | 0.2800 |
| Br | 3.013 | 0.4000 | 0.3000 | 0.3000 |
| 1 | 5.415 | 0.3600 | 0.3200 | 0.3200 |
| CTE | 1.061 | 0.4000 | 0.3000 | 0.3000 |
| CTR | 1.352 | 0.4150 | 0.2925 | 0.2925 |
| CBR | 1.896 | 0.3333 | 0.3333 | 0.3333 |
| CDI | 1.283 | 0.4500 | 0.2750 | 0.2750 |
| NTE | 0.964 | 0.4900 | 0.2550 | 0.2550 |
| NTR 2 | 1.030 | 0.4860 | 0.2570 | 0.2570 |
| $\mathrm{NPP12}$ | 1.090 | 0.5000 | 0.2500 | 0.2500 |
| ND1 | 0.956 | 0.4500 | 0.2750 | 0.2750 |
| OTE | 0.637 | 0.5200 | 0.2400 | 0.2400 |
| OTR4 | 0.569 | 0.4860 | 0.2570 | 0.2570 |
| OP12 | 0.274 | 0.4860 | 0.2570 | 0.2570 |
| STE | 3.000 | 0.4000 | 0.3000 | 0.3000 |
| STR4 | 3.729 | 0.4860 | 0.2570 | 0.2570 |
| SP12 | 2.700 | 0.4000 | 0.3000 | 0.3000 |
| PTE | 1.538 | 0.4000 | 0.3000 | 0.3000 |

${ }^{a}$ Atomic hybrids are defined in Table 1 of ref 10 . Scaled principal components of a tomic polarizability are related by $f_{\|}=f_{1}, f_{\perp}=f_{2}=f_{3}$ $=\left(1-f_{1}\right) / 2=\left(1-f_{\|}\right) / 2$ in the present paper.
(2) The dipole tensor is adjusted by omitting terms when $i$ is bonded to $j$, and the dipole interaction, eq 2 , is damped to yield
$T_{p i q j}=\left[-1 / r_{i j}{ }^{5}\right]\left[3 p_{i} q_{j}-\delta_{p q} r_{i j}{ }^{2}\right]\left[1-\exp \left[-s\left(r_{i j} / t \rho_{i j}\right)^{n}\right]\right]$
when $i$ is not bonded to $j$. The parameters $s=1.0, t=0.7$, and $n=10$ were found to be most effective.
(3) After the molecular polarizability is computed with these approximations to $\alpha_{i}(\mathrm{ahp})$ and $T_{\text {piqi; }}, \alpha_{\text {mol }}$ is scaled by the empirical molecular polarizability ${ }^{10}$ to yield the empirical molecular hybrid polarizability matrix

$$
\begin{equation*}
\alpha(\mathrm{ahp})=[\alpha(\mathrm{ahp}) / \alpha(\mathrm{mol})] \alpha_{\mathrm{mol}} \tag{17}
\end{equation*}
$$

where the average unscaled molecular polarizability satisfies

$$
\begin{equation*}
\operatorname{Tr}\left(\alpha_{\mathrm{mol}}\right) / 3=\alpha(\mathrm{mol}) \tag{18}
\end{equation*}
$$

and the molecular hybrid polarizability matrix satisfies

$$
\begin{equation*}
\operatorname{Tr}[\alpha(\mathrm{ahp})] / 3=\alpha(\mathrm{ahp}) \tag{19}
\end{equation*}
$$

The results are invariant to transformations of the entire molecule. The ahp, $\alpha$ (ahp), scaled principal components, $f_{1}, f_{2}, f_{3}$, unscaled molecular polarizabilities, $\alpha(\mathrm{mol})$, and orientation of the principal axes along with experimental results are presented in Table II. The rms deviation, between the empirical and experimental results is given by

$$
\begin{equation*}
\delta\left(f_{M}\right)=\left[\sum_{\mu}\left[f_{M_{\mu}}(\mathrm{ahp})-f_{M_{\mu}}(\mathrm{exp})\right]^{2}\right]^{1 / 2} \tag{20}
\end{equation*}
$$

with the scaled components $f_{M \mu}, \mu=1,2,3$, for each molecule $M$. The principal components of molecular polarizability are given by

$$
\begin{equation*}
\alpha_{M_{\mu}}(\text { calc })=\alpha(\text { calc }) f_{M_{\mu}}(\text { calc }) \quad \mu=1,2,3 \tag{21}
\end{equation*}
$$

where calc is ahp or exp. In the present method, the scaled polarizability tensor is adjusted to fit the scaled experimental values, while in the previous paper ${ }^{1}$ the average polarizability was parameterized to reproduce the average experimental values. They arc combined to yield $\alpha(\mathrm{ahp})$. Each approximation, is designed to reproduce essential parts of the polarizability tensor: the principal components and the average molecular polarizability while retaining atomic polarizabilities from earlier parameterizations.

The standard deviation in the calculation of principal components for $T$ molecules is the rms value

$$
\begin{equation*}
\sigma=\left[\sum_{M(=1 . T)} \delta\left(f_{M}\right)^{2} /(T-1)\right]^{1 / 2} \tag{22}
\end{equation*}
$$

In Table II the molecules are grouped as follows: diatomic and triatomic molecules, alkanes and halogenated alkanes, alkenes, benzene and its derivatives, alkynes, alcohols, ethers, aldehydes and ketones, nitriles, $\mathrm{NH}_{3}$ and amides, phosphates, DNA bases, rings containing $\mathrm{O}, \mathrm{N}, \mathrm{S}$, and condensed ring hydrocarbons. The molecules were built by using standard bond lengths and bond angles. The torsional angles are for minimum-energy conformations unless specified. In several cases polarizabilities for alternate conformations are presented.

To test the effect variations in the geometry has on $\alpha$ (ahp), the calculations were repeated by scaling all coordinates uniformly by 1.04 on 1.08 . The standard deviation changed by $0.0 \%$ and $0.1 \%$, respectively, for the individual molecules. This weak dependence on geometry is attributed to approximations I and 3; the constraints that $\alpha_{i}(\mathrm{ahp})$ are used for atoms and the scaling of the molecular polarizability matrix to $\alpha(\mathrm{ahp})$.

The results of the present theoretical method are compared to the empirical approaches of Applequist et al., ${ }^{7}$ Birge, ${ }^{8}$ and Thole ${ }^{9}$ in Table 111 only for cases in which experimental values are available. For most molecules, the present method yields results that are as close to experimental values as the best of these three theoretical methods. It was noted by Miller ${ }^{10}$ that the error in the average molecular polarizability is $0-4 \%$ for most molecules and in a few cases it may be as high as $8 \%$. Birge ${ }^{8}$ and Thole ${ }^{9}$ estimated the experimental error in the polarizability at $3 \%$ and in the components, as high as $6 \%$. Therefore, eqs 10, 16 , and 17 yield a good estimate of molecular polarizabilities. The important adjustment in the dipole tensor entails omitting the effect of bonded atoms and attenuating the contributions from nonbonded atoms. Use of anisotropic rather than isotropic atomic polarizabilities yields only an additional $1 \%$ improvement in $\alpha_{\text {mol }}$.

The principal axes are defined relative to symmetry planes and axes. For asymmetric molecules, the angles between the $\alpha_{1}, \alpha_{2}$, $\alpha_{3}$ principal axes and the bonds $X-Y$ are denoted by ( $\alpha_{1}, X Y$ ) etc. The conformations selected were minimum-energy conformations unless specified. When there are small barriers to rotation, polarizabilities of several low-energy conformations rather then the average are reported to demonstrate the dependence on conformation. In general, the dependence of polarizability on the conformation is small for the molecules studied. Torsion angles about CTE-CTE bonds are $\mathrm{t}, \mathrm{g}+$, and $\mathrm{g}-$, about $\mathrm{H}-\mathrm{CTE}-\mathrm{CTR}-$ OTR4 bonds are cis, and about H-NPI2-CTR-OTR4 bonds are trans. Methanol is cylindrically symmetrical with respect to rotation of the methyl (or OH ) group and the components of polarizability in the H-CTE-OTE-H conformations of $\mathrm{t}, \mathrm{g}+$, and $\mathrm{g}-$ remain unchanged. The cis and trans CTE-CTE-OTE-H conformations of ethanol average to $\left(f_{1}, f_{2}, f_{3}\right)=(0.2906,0.3291$, 0.3803 ) with an rms deviation of $0.6 \%$. The case of di-n-propyl ether is presented in Table IV. In the extended form of the backbone

the conformations are all trans. The van der Waals (6-14) potential energy, $U$, and the torsional energy, $V_{T}$, with a barrier of $3.0 \mathrm{kcal} \mathrm{mol}^{-1}$, about the CTE-CTE bonds are reported for $\mathrm{t}, \mathrm{g}+$, and $\mathbf{g}$ - forms for rotations about selective bonds. If only these conformations with $V_{\mathrm{T}}=0$ are considered in the weight averaging, then the rms error tends toward $5 \%$. The average orientation of the axes includes the mirror images, and consequently, $w$ possibilities occur for each state listed.

In the present approach, several considerations influenced the development of the empirical formulas. Their importance may be understood now by noting the results, presented in Table $V$, for several representative molecules with four successive ap-
proximations to $\alpha_{i}$ and $T_{i j}$ on seven representative molecules.
(1) If all $T_{\text {piqj }}=0$ and $\alpha_{i}$ are isotropic, then the three diagonal components of $f$ are equal and the rms error is 0.15 .
(2) If all $T_{p i q j}=0$ and anisotropic $\alpha_{i}$ are used, then f yields an rms error of 0.080 .
(3) If $T_{p i q j}$ is approximated by eq 16 and isotropic $\alpha_{i}$ are used, then the rms error is 0.065 .
(4) If $T_{p i q}$ and $\alpha_{i}$ are approximated by eqs 16 and 10 , then f reproduces the experimental polarizability components to an rms value of 0.035 .
$\mathrm{T}_{\text {piqj }}$ influences the computed unscaled molecular polarizabilities. Applequist et al. ${ }^{7}$ and Birge ${ }^{8}$ require smaller atomic polarizabilities, whereas Thole ${ }^{9}$ obtains some larger values then the additivity models of Vogel ${ }^{11}$ and the atomic hybrid method of Miller. ${ }^{10}$ However, if $\alpha_{\text {mol }}$ is scaled by the empirical value $\alpha$ (ahp), which reproduces the experimental values accurately, a good approximation is obtained to the components of molecular polarizability, and advantage is taken of the calibration by the empirical ahc and ahp methods of ref 10 . In approximation $\mathrm{I}, \mathrm{T}=0$ and isotropic atomic polarizabilities yield spherically symmetrical molecular polarizabilities. The principal axes are arbitrary. In the second approximation, $\mathrm{T}=0$ and anisotropic atomic polarizabilities yield improved results for the standard deviation, but the principal axes do not agree with the experimental results as well as with nonzero values of $T_{\text {pief: }}$. Approximation $3, \mathrm{~T} \neq 0$ and isotropic atomic polarizabilities also yield an improved standard deviation $\sigma$, but the principal axes do not consistently agree with experimental results. Approximation 4 combines the essential features, and the best results are obtained when $T \neq 0$ and anisotropic atomic polarizabilities are used according to eqs 10 , 16, and 17. Both are required to minimize the standard deviation and direct the principal axes in agreement with experimental results.

## Discussion and Conclusion

The components of molecular polarizability are computed with approximations 10,16 , and 17 . The optimum scaled components of atomic polarizability are reported in Table I along with the average atomic polarizability, $\alpha_{A}(a h p)$, from ref 10 . The components add to 1 , and the perpendicular components are set equal so that each contribution to $\alpha_{i}$ is cylindrically symmetrical about the direction $i$ to $j$ to simplify the transformation from the local coordinate system directed along $i$ to $j$ to the space-fixed system. They were initially chosen to reproduce the components of molecular polarizability in diatomic molecules and then adjusted to fit and optimize the entire set of data. The parameters $s=1.00$, $t=0.700$, and $n=10$ are found to phase in the dipole approximation so that $T_{\text {piqj }}$ approaches the classical dipole tensor, eq 2 , as $r_{i j}$ exceeds $t \rho_{i j}=0.7 \rho_{i j}$, which is approximately $1.5-2$ bond lengths. The point dipole approximation is valid at large distances but must be modified at distances less than the sum $\rho_{i j}$. Thole ${ }^{9}$ discussed this in terms of the atomic polarizabilities of atoms intersecting each other. In the present approximation, the effect of bonded atoms is included through anisotropic atomic polariz-

[^1]Table II. Polarizabilities, Scaled Components, and Orientation of Principal Axes

| molecule | $\alpha,{ }^{a} \AA^{3}$ | $f_{1}{ }^{\text {b }}$ | $f_{2}{ }^{\text {b }}$ | $f_{3}{ }^{6}$ | $\alpha(\mathrm{mol}),{ }^{\text {c }} \AA^{3}$ | \% $\mathrm{rms}^{\text {d }} \sigma(f)$ |  | refe | orientation |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{H}_{2}$ | 0.77 | 0.3030 | 0.3030 | 0.3940 | 0.77 | 0.2 | M |  | $\alpha_{3} \\|$ symmetry axis |
|  | 0.79 | 0.3038 | 0.3038 | 0.3924 |  |  | E | LB |  |
| $\mathrm{N}_{2}$ | 1.91 | 0.2750 | 0.2750 | 0.4500 | 1.91 | 0.1 | M |  | $\alpha_{3} \\|$ symmetry axis |
|  | 1.76 | 0.2746 | 0.2746 | 0.4508 |  |  | E | LB |  |
| $\mathrm{O}_{2}$ | 1.14 | 0.2570 | 0.2570 | 0.4860 | 1.14 | 0.8 | M |  | $\alpha_{3} \\|$ symmetry axis |
|  | 1.60 | 0.2536 | 0.2536 | 0.4926 |  |  | E | LB |  |
| $\mathrm{Cl}_{2}$ | 4.63 | 0.2800 | 0.2800 | 0.4400 | 4.63 | 4.5 | M |  | $\alpha_{3} \\|$ symmetry axis |
|  | 4.61 | 0.2617 | 0.2617 | 0.4772 |  |  | E | LB |  |
| HF | 0.68 | 0.2889 | 0.2889 | 0.4223 | 0.68 | 2.7 | M |  | $\alpha_{3} \\|$ symmetry axis |
|  | 0.80 | 0.3000 | 0.3000 | 0.4000 |  |  | E | LB |  |
| HCl | 2.70 | 0.2833 | 0.2833 | 0.4334 | 2.70 | 4.6 | M |  | $\alpha_{3} \\|$ symmetry axis |
|  | 2.63 | 0.3029 | 0.3029 | 0.3967 |  |  | E | LB |  |
| HBr | 3.40 | 0.3003 | 0.3003 | 0.3993 | 3.40 | 1.2 | M |  | $\alpha_{3} \\|$ symmetry axis |
|  | 3.61 | 0.3056 | 0.3056 | 0.3897 |  |  | E | LB |  |
| HI | 5.80 | 0.3189 | 0.3189 | 0.3623 | 5.80 | 4.9 | M |  | $\alpha_{3} \\|$ symmetry axis |
|  | 5.45 | 0.2991 | 0.2991 | 0.4024 |  |  | E | LB |  |
| CO | 1.92 | 0.2820 | 0.2820 | 0.4360 | 1.92 | 1.0 | M |  | $\alpha_{3} \\|$ symmetry axis |
|  | 1.95 | 0.2778 | 0.2778 | 0.4444 |  |  | E | LB |  |
| $\mathrm{CO}_{2}$ | 2.42 | 0.2535 | 0.2535 | 0.4930 | 2.48 | 1.4 | M |  | $\alpha_{3} \\|$ symmetry axis |
|  | 2.65 | 0.2478 | 0.2478 | 0.5044 |  |  | E | LB |  |
| $\mathrm{SO}_{2}$ | 5.00 | 0.2456 | 0.3044 | 0.4500 | 5.06 | 2.4 | M |  | $\alpha_{1} \perp$ OSO plane, $\alpha_{2} \\|$ 2-fold axis |
|  | 3.90 | 0.2325 | 0.2983 | 0.4692 |  |  | E | LB |  |
| $\mathrm{CS}_{2}$ | 8.74 | 0.2026 | 0.2026 | 0.5948 | 10.06 | 2.1 | M |  | $\alpha_{3} \\|$ symmetry axis |
|  | 8.74 | 0.2113 | 0.2113 | 0.5774 |  |  | E | LB |  |
| $\mathrm{H}_{2} \mathrm{~S}$ | 3.77 | 0.2959 | 0.3420 | 0.3621 | 3.79 | 1.3 | M |  | $\alpha_{1} \perp$ plane, $\alpha_{2} \\| 2$-fold axis |
|  | 3.83 | 0.2994 | 0.3490 | 0.3516 |  |  | E | LB |  |
| $\mathrm{N}_{2} \mathrm{O}$ | 2.48 | 0.2466 | 0.2466 | 0.5068 | 2.60 | 4.1 | M |  | $\alpha_{3} \\|$ symmetry axis |
|  | 3.00 | 0.2300 | 0.2300 | 0.5400 |  |  | E | LB |  |
| methane | 2.61 | 0.3333 | 0.3333 | 0.3333 | 2.68 | 0.0 | M |  | regular tetrahedron |
|  | 2.60 | 0.3333 | 0.3333 | 0.3333 |  |  | E | LB |  |
| ethane | 4.44 | 0.3142 | 0.3142 | 0.3715 | 4.64 | 0.1 | M |  | $\alpha_{3} \\| 3$-fold axis |
|  | 4.48 | 0.3140 | 0.3140 | 0.3720 |  |  | E | AP |  |
| propane | 6.28 | 0.2993 | 0.3197 | 0.3810 | 6.60 | 2.8 | M |  | $\alpha_{1} \perp$ CCC plane, $\alpha_{2} \\|$ 2-fold axis |
|  | 6.38 | 0.2999 | 0.2999 | 0.4002 |  |  | E | AP |  |
| neopentane | 9.95 | 0.3333 | 0.3333 | 0.3333 | 10.48 | 0.0 | M |  | regular tetrahedron |
|  | 10.20 | 0.3333 | 0.3333 | 0.3333 |  |  | E | AP |  |
| cyclopentane | 9.18 | 0.2890 | 0.3533 | 0.3577 | 9.77 | 2.6 | M |  | $\alpha_{1} \perp$ average plane |
|  | 9.15 | 0.3060 | 0.3341 | 0.3599 |  |  | E | AP |  |
| cyclopentane | 9.18 | 0.2890 | 0.3533 | 0.3577 | 9.77 | 2.6 | M |  | $\alpha_{1} \perp$ average plane |
|  | 9.15 | 0.3060 | 0.3341 | 0.3599 |  |  | E | AP |  |
| cyclohexane | 11.01 | 0.2841 | 0.3580 | 0.3580 | 11.67 | 0.1 | M |  | $\alpha_{1} \perp 3$-fold axis and \|| to axial |
|  | 11.00 | 0.2842 | 0.3579 | 0.3579 |  |  | E | AP | CH bonds |
| $\mathrm{CH}_{3} \mathrm{~F}$ | 2.52 | 0.3316 | 0.3316 | 0.3369 | 2.58 | 8.1 | M |  | $\alpha_{3} \\| 3$-fold axis |
|  | 2.62 | 0.2977 | 0.2977 | 0.4046 |  |  | E | AP |  |
| $\mathrm{CHF}_{3}$ | 2.34 | 0.3298 | 0.3351 | 0.3351 | 2.39 | 1.3 | $\mathrm{M}$ |  | $\alpha_{1} \\| 3$-fold axis |
|  | 2.81 | 0.3191 | 0.3405 | 0.3405 |  |  | E | AP |  |
| $\mathrm{CF}_{4}$ | 2.25 | 0.3333 | 0.3333 | 0.3333 | 2.26 | 0.0 | M |  | regular tetrahedron |
|  | 2.92 | 0.3333 | 0.3333 | 0.3333 |  |  | E | AP |  |
| $\mathrm{CH}_{3} \mathrm{Cl}$ | 4.54 | 0.2848 | 0.2848 | 0.4305 | 4.76 | 1.7 | M |  | $\alpha_{3}\| \| 3$-fold axis |
|  | 4.55 | 0.2916 | 0.2916 | 0.4161 |  |  | E | AP |  |
| $\mathrm{CH}_{2} \mathrm{Cl}_{2}$ | 6.47 | 0.2563 | 0.3124 | 0.4313 | 6.96 | 0.7 | M |  | $\alpha_{1} \perp \mathrm{ClCCl}$ plane, $\alpha_{2} \\|$ 2-fold axis |
|  | 6.82 | 0.2620 | 0.3079 | 0.4306 |  |  | E | AP |  |
| $\mathrm{CHCl}_{3}$ | 8.39 | 0.2542 | 0.3729 | 0.3729 | 9.25 | 1.1 | M |  | $\alpha_{1} \\| 3$-fold axis |
|  | 8.53 | 0.2634 | 0.3681 | 0.3681 |  |  | E | AP |  |
| $\mathrm{CCl}_{4}$ | 10.32 | 0.3333 | 0.3333 | 0.3333 | 11.58 | 0.0 | M |  | regular tetrahedron |
|  | 10.51 | 0.3333 | 0.3333 | 0.3333 |  |  | E | AP |  |
| $\mathrm{C}_{2} \mathrm{H}_{5} \mathrm{Cl}$ | 6.37 | 0.2770 | 0.3086 | 0.4143 | 6.79 | 4.3 | $\mathbf{M}$ |  |  |
|  | 6.40 | 0.2609 | 0.3438 | 0.3953 |  |  | E | LB | $\left(\alpha_{2}, \mathrm{CCl}\right)=69.9^{\circ}$ |
| $\mathrm{CH}_{3} \mathrm{Br}$ | 5.24 | 0.2940 | 0.2940 | 0.4121 | 5.44 | 0.2 | $\mathrm{M}$ |  | $\alpha_{3}$ \|| 3 -fold axis |
|  | 5.61 | 0.2947 | 0.2947 | 0.4106 |  |  | E | AP | $\alpha_{3} \mid$ |
| $\mathrm{CH}_{2} \mathrm{Br}_{2}$ | 7.86 | 0.2754 | 0.2861 | 0.4385 | 8.30 |  | M |  | $\alpha_{1} \perp \operatorname{BrCBr}$ plane, $\alpha_{2} \\|$ 2-fold axis |
|  | 8.68 | 0.0000 | 0.0000 | 0.0000 |  |  | E | AP |  |
| $\mathrm{CHBr}_{3}$ | 10.49 | 0.2624 | 0.3688 | 0.3688 | 11.31 | 0.7 | M |  | $\alpha_{1}\| \| 3$-fold axis |
|  | 11.84 | 0.2683 | 0.3660 | 0.3660 |  |  | E | AP |  |
| $\mathrm{CH}_{3} \mathrm{l}$ | 7.64 | 0.3046 | 0.3046 | 0.3907 | 7.86 | 0.7 | M |  | $\alpha_{3}\| \| 3$-fold axis |
|  | 7.59 | 0.3017 | 0.3017 | 0.3961 |  |  | E | AP |  |
| $\mathrm{CH}_{2} \mathrm{I}_{2}$ | 12.67 12.90 | 0.2782 | 0.3050 | 0.4168 | 13.27 |  | M |  | $\alpha_{1} \perp$ ICI plane, $\alpha_{2}\| \| 2$-fold axis |
|  | 12.90 | 0.0000 | 0.0000 | 0.0000 |  |  | E | AP |  |
| $\mathrm{CHI}_{3}$ | 17.69 18.04 | 0.2627 0.3093 | 0.3686 | 0.3686 | 18.83 | 5.7 | M |  | $\alpha_{1}\| \| 3$-fold axis |
|  | 18.04 | 0.3093 | 0.3453 | 0.3453 |  |  | E | AP |  |
| ethylene | 4.25 | 0.2509 | 0.3509 | 0.3982 | 4.44 | 8.6 | M |  | $\alpha_{1} \perp$ plane, $\alpha_{3} \\| \mathrm{C}=\mathrm{C}$ axis |
|  | 4.26 | 0.2809 | 0.2809 | 0.4390 |  |  | E | LB |  |
| ethylene | 4.25 | 0.2528 | 0.3516 | 0.3956 | 4.43 |  | M |  | $\alpha_{1} \perp$ plane, $\alpha_{3} \\| \mathrm{C}=\mathrm{C}$ axis |
|  | 4.26 | 0.0000 | 0.0000 | 0.0000 |  |  | E | BR |  |
| $\mathrm{CH}_{2}=\mathrm{CCl}_{2}$ | 8.11 | 0.2237 | 0.3643 | 0.4120 | 8.86 | 4.0 | M |  | $\alpha_{1} \perp$ plane, $\alpha_{2} \\| \mathrm{C}=\mathrm{C}$ axis |
|  | 7.83 | 0.2448 | 0.3742 | 0.3814 |  |  | E | BR |  |

Table II (Continued)


Table II (Continued)

| molecule | $\alpha,{ }^{6} \AA^{3}$ | $f_{1}{ }^{\text {b }}$ | $f_{2}{ }^{\text {b }}$ | $f_{3}{ }^{6}$ | $\alpha(\mathrm{mol}),{ }^{\text {c }}{ }^{\text {A }}{ }^{3}$ | $\begin{gathered} \% \mathrm{rms}^{d} \\ \sigma(f) \end{gathered}$ |  | ref ${ }^{\text {e }}$ | orientation |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| formaldehyde | 2.70 | 0.2636 | 0.3416 | 0.3948 | 2.78 | 4.2 | M |  | $\alpha_{1} \perp$ plane, $\alpha_{3} \\| 2$-fold axis |
|  | 2.45 | 0.2490 | 0.3755 | 0.3755 |  |  | E | AP |  |
| acetaldehyde | 4.53 | 0.2782 | 0.3361 | 0.3857 | 4.76 |  | M |  | $\alpha_{1} \perp \mathrm{C}(\mathrm{C}=\mathrm{O}) \mathrm{C}$ plane, $\left(\alpha_{2}, \mathrm{CC}\right)=110^{\circ}$, |
|  | 4.590 | 0.000 | 0.0000 | 0.0000 |  |  | E | AP | $\begin{aligned} & \left(\alpha_{2}, \mathrm{C}=\mathrm{O}\right)=130^{\circ},\left(\alpha_{3}, \mathrm{CC}\right)=20^{\circ}, \\ & \left(\alpha_{3}, \mathrm{C}=0\right)=140^{\circ} \end{aligned}$ |
| acetone | 6.37 | 0.2768 | 0.3425 | 0.3807 | 6.74 | 6.3 | M |  | $\alpha_{1} \perp \mathrm{C}(\mathrm{C}=0) \mathrm{C}$ plane, $\alpha_{2} \\| \mathrm{C}=02$-fold axis |
|  | 6.39 | 0.2306 | 0.3845 | 0.3845 |  |  | E | AP |  |
| acetone | 6.37 | 0.2768 | 0.3425 | 0.3807 | 6.74 | 3.8 | M |  | $\alpha_{1} \perp \mathrm{C}(\mathrm{C}=0) \mathrm{C}$ plane, $\alpha_{2} \\| \mathrm{C}=0$ 2-fold axis |
|  | 6.39 | 0.2546 | 0.3725 | 0.3735 |  |  | E | AP |  |
| methyl ethyl ketone | 8.20 | 0.2712 | 0.3217 | 0.4071 | 8.79 | 3.8 | M |  | $\alpha_{1} \perp \mathrm{C}(\mathrm{C}=\mathrm{O}) \mathrm{C}$ plane, $\left(\alpha_{2}, \mathrm{C}=\mathrm{O}\right)=14^{\circ}$, |
|  | 8.13 | 0.2468 | 0.3506 | 0.4030 |  |  | E | LB | $\left(\alpha_{3}, \mathrm{C}=\mathrm{O}\right)=104^{\circ}, C_{s}$ symmetry |
| diethyl ketone | 10.04 | 0.2662 | 0.3137 | 0.4202 | 10.88 | 3.5 | M |  | $\alpha_{1} \perp \mathrm{C}(\mathrm{C}=0) \mathrm{C}$ plane, $\alpha_{2} \\| \mathrm{C}=0$ |
|  | 9.93 | 0.2400 | 0.3360 | 0.4243 |  |  | E | LB |  |
| methyl propyl ketone | 10.04 | 0.2660 | 0.3115 | 0.4225 | 10.87 | 3.6 | M |  | $\alpha_{1} \perp \mathrm{C}(\mathrm{C}=\mathrm{O}) \mathrm{C}$ plane, $\left(\alpha_{2}, \mathrm{C}=\mathrm{O}\right)=3^{\circ}$, |
|  | 9.93 | 0.2551 | 0.3404 | 0.4048 |  |  | E | LB | $\left(\alpha_{3}, \mathrm{C}=\mathrm{O}\right)=87^{\circ}, C_{s}$ symmetry |
| diisopropyl ketone | 13.71 | 0.2938 | 0.3309 | 0.3753 | 14.62 | 3.1 | M |  | $\alpha_{2} \perp \mathrm{C}(\mathrm{C}=\mathrm{O}) \mathrm{C} \text { plane, }\left(\alpha_{1}, \mathrm{C}=\mathrm{O}\right)=161^{\circ}$ |
|  | 13.53 | 0.2882 | 0.3117 | 0.3991 |  |  | E | LB | $\left(\alpha_{2}, C=0\right)=109^{\circ}, C_{s}$ symmetry |
| diisopropyl ketone | 13.71 | 0.2832 | 0.3305 | 0.3863 | 14.71 | 2.3 | M |  | $\alpha_{2} \perp \perp \mathrm{C}(\mathrm{C}=\mathrm{O}) \mathrm{C}$ plane, $\alpha_{1} \\| \mathrm{C}=\mathrm{O}$ axis, |
|  | 13.53 | 0.2882 | 0.3117 | 0.3991 |  |  | E | LB | $C_{2 v}$ symmetry |
| HCN | 2.63 | 0.2643 | 0.2643 | 0.4715 | 2.68 | 4.1 | M |  | $\alpha_{3}$ \|| symmetry axis |
|  | 2.59 | 0.2471 | 0.2471 | 0.5045 |  |  | E | LB |  |
| $(\mathrm{CN})_{2}$ | 4.48 | 0.2261 | 0.2261 | 0.5477 | 4.94 | 3.9 | M |  | $\alpha_{3}\| \|$ symmetry axis |
|  | 5.01 | 0.2422 | 0.2422 | 0.5163 |  |  | E | LB |  |
| acetonitrile | 4.46 | 0.2758 | 0.2758 | 0.4484 | 4.74 | 2.6 | M |  | $\alpha_{3}\| \|$ symmetry axis |
|  | 4.48 | 0.2865 | 0.2865 | 0.4271 |  |  | E | AP |  |
| $\left(\mathrm{CH}_{3}\right)_{3} \mathrm{CCN}$ | 9.97 | 0.3168 | 0.3168 | 0.3665 | 10.70 | 0.7 | M |  | $\alpha_{3}\| \|$ symmetry axis |
|  | 9.59 | 0.3139 | 0.3139 | 0.3723 |  |  | E | AP |  |
| $\mathrm{CCl}_{3} \mathrm{CN}$ | 10.25 | 0.3299 | 0.3299 | 0.3402 | 11.54 | 0.2 | M |  | $\alpha_{3}\| \|$ symmetry axis |
|  | 10.42 | 0.3292 | 0.3292 | 0.3420 |  |  | E | AP |  |
| $\mathrm{NH}_{3}$ | 2.13 | 0.2761 | 0.3619 | 0.3619 | 2.16 | 6.1 | $\mathbf{M}$ |  | $\alpha_{1} \perp$ to 3 -fold axis |
|  | 2.26 | 0.3215 | 0.3215 | 0.3569 |  |  | E | LB |  |
| $\mathrm{HCONH}_{2}$ | 4.17 | 0.2317 | 0.3573 | 0.4110 | 4.40 | 2.4 | $\mathrm{M}$ |  |  |
|  | 4.08 | 0.0000 | 0.5727 | 0.4281 |  |  | E | AP | $\begin{aligned} & \left(\alpha_{2}, \mathrm{CN}\right)=113^{\circ},\left(\alpha_{3}, \mathrm{C}=\mathrm{O}\right)=37^{\circ}, \\ & \left(\alpha_{3}, \mathrm{CN}\right)=157^{\circ} \end{aligned}$ |
| $\mathrm{HCONHCH}_{3}$ | 6.01 | 0.2456 | 0.3284 | 0.4260 | 6.45 | 6.2 | M |  | $\alpha_{1} \perp$ plane, $\left(\alpha_{2}, \mathrm{C}=\mathrm{O}\right)=111^{\circ}$, |
|  | 5.91 | 0.0000 | 0.6176 | 0.3818 |  |  | E | AP | $\begin{aligned} & \left(\alpha_{2}, \mathrm{CN}\right)=129^{\circ},\left(\alpha_{3}, \mathrm{C}=\mathrm{O}\right)=21^{\circ} \\ & \left(\alpha_{3}, \mathrm{CN}\right)=141^{\circ}, \mathrm{CH}_{3} \text { trans to } \mathrm{C}=\mathrm{O} \end{aligned}$ |
| $\mathrm{HCONHCH}_{3}$ | 6.01 | 0.2465 | 0.3503 | 0.4032 | 6.41 | 3.0 | M |  | $\alpha_{1} \perp$ plane, $\left(\alpha_{2}, \mathrm{C}=\mathrm{O}\right)=1^{\circ}$, |
|  | 5.91 | 0.0000 | 0.6176 | 0.3818 |  |  | E | AP | $\begin{aligned} & \left(\alpha_{2}, \mathrm{CN}\right)=119^{\circ},\left(\alpha_{3}, \mathrm{C}=\mathrm{O}\right)=89^{\circ}, \\ & \left(\alpha_{3}, \mathrm{CN}\right)=151^{\circ}, \mathrm{CH} 3 \text { cis to } \mathrm{C}=\mathrm{O} \end{aligned}$ |
| $\mathrm{HCON}\left(\mathrm{CH}_{3}\right)_{2}$ | 7.84 | 0.2486 | 0.3648 | 0.3866 | 8.44 | 2.4 | M |  | $\alpha_{1} \perp$ plane, $\left(\alpha_{2}, \mathrm{C}=\mathrm{O}\right)=92^{\circ}$, |
|  | 7.81 | 0.0000 | 0.5962 | 0.4033 |  |  | E | AP | $\begin{aligned} & \left(\alpha_{2}, \mathrm{CN}\right)=148^{\circ},\left(\alpha_{3}, \mathrm{C}=\mathrm{O}\right)=2^{\circ}, \\ & \left(\alpha_{3}, \mathrm{CN}\right)=122^{\circ} \end{aligned}$ |
| $\mathrm{CH}_{3} \mathrm{CONH}_{2}$ | $6.01$ | $0.2485$ | $0.3559$ | $0.3956$ | 6.41 | 0.2 | $\mathrm{M}$ |  |  |
|  | 5.67 | $0.0000$ | $0.6055$ | $0.3939$ |  |  | E | AP | $\begin{aligned} & \left(\alpha_{2}, \mathrm{CN}\right)=56^{\circ},\left(\alpha_{3}, \mathrm{C}=\mathrm{O}\right)=94^{\circ}, \\ & \left(\alpha_{3}, \mathrm{CN}\right)=146^{\circ} \end{aligned}$ |
| $\mathrm{CH}_{3} \mathrm{C}\left(\mathrm{CH}_{2} \mathrm{O}\right)_{3} \mathrm{P}$ | 12.24 | 0.3097 | 0.3097 | 0.3805 | 12.96 | 1.9 | M |  | $\alpha_{3}\| \| 3$-fold axis |
|  | 13.32 | 0.3173 | 0.3173 | 0.3654 |  |  | E | AL |  |
| $\mathrm{CH}_{3} \mathrm{C}\left(\mathrm{CH}_{2} \mathrm{O}\right)_{3} \mathrm{PO}$ | 12.87 | 0.2978 | 0.2980 | 0.4042 | 13.83 | 2.9 | M |  | $\alpha_{3}\| \| 3$-fold axis |
|  | 12.84 | 0.3097 | 0.3097 | 0.3806 |  |  | E | AL |  |
| $\mathrm{CH}_{3} \mathrm{C}\left(\mathrm{CH}_{2} \mathrm{O}\right)_{3} \mathrm{PS}$ | 15.24 | 0.2855 | 0.2859 | 0.4286 | 16.57 | 2.7 | M |  | $\alpha_{3}\| \| 3$-fold axis |
|  | 15.74 | 0.2967 | 0.2967 | 0.4064 |  |  | E | AL |  |
| uracil | 10.27 | 0.1888 | 0.3918 | 0.4194 | 11.81 | 5.7 | M |  | $\alpha_{1} \perp$ plane, $\left(\alpha_{2}, \mathrm{~N} 3 \mathrm{H} 3\right)=5^{\circ}$, |
|  | 10.09 | 0.1943 | 0.3492 | 0.4569 |  |  | E | BG | $\begin{aligned} & \left(\alpha_{3}, \mathrm{~N} 3 \mathrm{H} 3\right)=95^{\circ},\left(\alpha_{2}, \mathrm{C} 4 \mathrm{O} 4\right)=52^{\circ}, \\ & \left(\alpha_{3}, \mathrm{C} 4 \mathrm{O} 4\right)=38^{\circ} \end{aligned}$ |
| guanine | 15.68 | 0.1846 | 0.3570 | 0.4585 | 18.54 | 3.2 | M |  |  |
|  | 13.60 | 0.1799 | 0.3819 | 0.4385 |  |  | E | BO | $\begin{aligned} & \left(\alpha_{3}, \mathrm{C} 2 \mathrm{~N} 2\right)=20^{\circ},\left(\alpha_{2}, \mathrm{C} 6 \mathrm{O} 6\right)=0^{\circ}, \\ & \left(\alpha_{3}, \mathrm{C} 6 \mathrm{O} 6\right)=90^{\circ} \end{aligned}$ |
| adenine | 15.05 | 0.1911 | 0.3864 | 0.4225 | 17.40 | 0.8 | M |  | $\alpha_{1} \perp$ plane, ( $\left.\alpha_{2}, \mathrm{C} 6 \mathrm{~N} 6\right)=29^{\circ}$, |
|  | 13.10 | 0.1863 | 0.3852 | 0.4282 |  |  | E | BO | $\begin{aligned} & \left(\alpha_{3}, \mathrm{C} 6 \mathrm{~N} 6\right)=61^{\circ},\left(\alpha_{2}, \mathrm{C} 8 \mathrm{H} 8\right)=66^{\circ}, \\ & \left(\alpha_{3}, \mathrm{C} 8 \mathrm{H} 8\right)=156^{\circ} \end{aligned}$ |
| cytosine | 11.12 | 0.1868 | 0.3679 | 0.4452 | 12.87 | 3.3 | M |  | $\alpha_{1} \perp$ plane, $\left(\alpha_{2}, \mathrm{C} 2 \mathrm{O} 2\right)=87^{\circ}$, |
|  | 10.30 | 0.1919 | 0.3421 | 0.4654 |  |  | E | BO | $\begin{aligned} & \left(\alpha_{3}, \mathrm{C} 2 \mathrm{O} 2\right)=177^{\circ},\left(\alpha_{2}, \mathrm{C} 4 \mathrm{~N} 4\right)=34^{\circ}, \\ & \left(\alpha_{3}, \mathrm{C} 4 \mathrm{~N} 4\right)=56^{\circ} \end{aligned}$ |
| thymine | 12.11 | 0.1950 | 0.3625 | 0.4425 | 14.03 | 1.1 | M |  | $\alpha_{1} \perp$ plane, $\left(\alpha_{2}, \mathrm{C} 4 \mathrm{O} 4\right)=72^{\circ}$, |
|  | 11.23 | 0.2009 | 0.3535 | 0.4455 |  |  | E | BO | $\begin{aligned} & \left(\alpha_{3}, \mathrm{C} 4 \mathrm{O} 4\right)=18^{\circ},\left(\alpha_{2}, \mathrm{C} 5 \mathrm{Me}\right)=133^{\circ}, \\ & \left(\alpha_{3}, \mathrm{C} 5 \mathrm{Me}\right)=43^{\circ} \end{aligned}$ |
| furan | 7.23 | 0.2290 | 0.3634 | 0.4076 | 7.78 | 2.8 | M |  | $\alpha_{1} \perp$ plane, $\alpha_{2}\| \| 2$-fold axis |
|  | 7.23 | 0.2476 | 0.3426 | 0.4099 |  |  | E | LL |  |
| thiophene | 9.66 | 0.2308 | 0.3675 | 0.4018 | 10.41 | 3.2 | M |  | $\alpha_{1} \perp$ plane, $\alpha_{2}\| \| 2$-fold axis |
|  | 9.00 | 0.2481 | 0.3756 | 0.3759 |  |  | E | LL |  |
| pyrrole | 8.43 | 0.2138 | 0.3868 | 0.3995 | 9.28 | 3.7 | M |  | $\alpha_{1} \perp$ plane, $\alpha_{3}\| \| 2$-fold axis |
|  | 7.94 | 0.2431 | 0.3648 | 0.3917 |  |  | E | LL |  |
| pyridine | 9.73 | 0.2029 | 0.3762 | 0.4209 | 10.84 | 0.6 | M |  | $\alpha_{1} \perp$ plane, $\alpha_{2}\| \| 2$-fold axis |
|  | 9.47 | 0.2034 | 0.3816 | 0.4182 |  |  | E | LB |  |
| pyridine | 10.11 | 0.2003 | 0.3974 | 0.4023 | 11.36 | 4.1 | M |  | $\alpha_{1} \perp$ plane, $\alpha_{2}\| \| 2$-fold axis |

Table II (Continued)

| molecule | $\alpha,{ }^{6} \AA^{3}$ | $f_{1}{ }^{\text {b }}$ | $f_{2}{ }^{\text {b }}$ | $f_{3}{ }^{\text {b }}$ | $\alpha(\mathrm{mol}),{ }^{\text {c }}{ }^{\text {d }}{ }^{3}$ | $\% \mathrm{rms}^{d} \sigma(f)$ |  | refe | orientation |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| naphthalene | 9.20 | 0.2337 | 0.3779 | 0.3884 | 20.93 | 5.6 | E | LL | $\alpha_{1} \perp$ plane, $\alpha_{3} \\|$ (minor) 2 -fold axis, $\alpha_{3} \\|$ (major) 2 -fold axis |
|  | 17.70 | 0.1822 | 0.3599 | 0.4579 |  |  | M |  |  |
|  | 16.57 | 0.2273 | 0.3339 | 0.4385 |  | 3.4 | E | LR |  |
| quinoline | 17.00 | 0.1836 | 0.3481 | 0.4684 | 20.08 |  | M | LL | $\begin{aligned} & \alpha_{1} \perp \mathrm{pl},\left(\alpha_{2}, \mathrm{~N} 1 \cdots \mathrm{C} 4\right)=2^{\circ},\left(\alpha_{3},\right. \\ & \mathrm{N} 1 \cdots \mathrm{C} 4)=88^{\circ},\left(\alpha_{2}, \mathrm{~N} 1 \cdots \mathrm{C} 8\right)=88^{\circ}, \\ & \left(\alpha_{3}, \mathrm{~N} 1 \cdots \mathrm{C} 8\right)=178^{\circ} \end{aligned}$ |
|  | 15.68 | 0.2017 | 0.3571 | 0.4409 |  |  |  |  |  |
| isoquinoline | 17.00 | 0.1845 | 0.3652 | 0.4503 | 19.90 | 5.4 | $\begin{aligned} & \mathrm{M} \\ & \mathrm{E} \end{aligned}$ | LL | $\begin{aligned} & \alpha_{1} \stackrel{\perp}{\perp} \mathrm{pl},\left(\alpha_{2}, \mathrm{~N} 2 \cdots \mathrm{C} 3\right)=9^{\circ},\left(\alpha_{3},\right. \\ & \mathrm{N} 2 \cdots \mathrm{C} 3)=81^{\circ},\left(\alpha_{2}, \mathrm{~N} 2 \cdots \mathrm{C} 7\right)=81^{\circ}, \\ & \left(\alpha_{3}, \mathrm{~N} 2 \cdots \mathrm{C} 7\right)=171^{\circ} \end{aligned}$ |
|  | 15.62 | 0.1660 | 0.4085 | 0.4247 |  |  |  |  |  |
| anthracene | 24.97 | 0.1663 | 0.3264 | 0.5072 | 31.22 | 5.4 | $\begin{aligned} & \mathrm{M} \\ & \mathrm{E} \end{aligned}$ |  | $\alpha_{1} \perp \text { plane, } \alpha_{2} \\| \text { (minor) 2-fold axis, }$ |
|  | 25.37 | 0.2037 | 0.3285 | 0.4677 |  |  |  | LR | $\alpha_{3} \\|$ (major) 2 -fold axis |
| pyrene | 28.77 | 0.1603 | 0.3778 | 0.4619 | 36.04 | 5.5 | $\begin{aligned} & \mathrm{M} \\ & \mathrm{E} \end{aligned}$ | LR | $\alpha_{1} \perp$ plane, $\alpha_{2} \\|$ (minor) 2 -fold axis, $\alpha_{3} \\|$ (major) 2 -fold axis |
|  | 28.22 | 0.2008 | 0.3402 | 0.4589 |  |  |  |  |  |
| chrysene | 32.24 | 0.1528 | 0.3211 | 0.5261 | 42.18 | 9.4 | $\begin{aligned} & \mathrm{M} \\ & \mathrm{E} \end{aligned}$ |  | $\begin{aligned} & \alpha_{1} \stackrel{\perp}{ } \mathrm{pl},\left(\alpha_{2}, \mathrm{C} 3 \cdots \mathrm{C} 2\right)=22^{\circ},\left(\alpha_{3},\right. \\ & \mathrm{C} 3 \cdots \mathrm{C} 2)=68^{\circ},\left(\alpha_{2}, \mathrm{C} 3 \cdots \mathrm{C} 9\right)=68^{\circ}, \\ & \left(\alpha_{3}, \mathrm{C} 3 \cdots \mathrm{C} 9\right)=158^{\circ} \end{aligned}$ |
|  | 32.15 | 0.2011 | 0.3406 | 0.4479 |  |  |  | LR |  |
| triphenylene | 32.24 | 0.1533 | 0.4233 | 0.4233 | 41.46 | 1.5 | $\begin{aligned} & \mathrm{M} \\ & \mathrm{E} \end{aligned}$ |  | $\alpha_{1} \perp$ plane |
|  | 31.07 | 0.1652 | 0.4173 | 0.4173 |  |  |  | LS |  |
| coronene | 43.62 | 0.1409 | 0.4295 | 0.4295 | 58.16 | 1.6 | M |  | $\alpha_{1} \perp$ plane |
|  | 44.77 | 0.1541 | 0.4229 | 0.4229 |  |  | E | LS |  |

${ }^{a} \alpha$, line M , refers to $\alpha(\mathrm{ahp})=\alpha(\mathrm{mhp}) ;{ }^{10} \alpha$, line E, refers to the average experimental result. ${ }^{b} f_{1}, f_{2}$, and $f_{3}$, the scaled principal components of polarizability, sum to 1 within round off and experimental error. ${ }^{c} \alpha$ (mol), the average unscaled molecular polarizability, eq 9 , is included to demonstrate the effect of scaling. ${ }^{d}$ The rms deviation is defined by eq 20. e'References: LB, ref 12; AP, ref 7; BR, ref 13; LS, ref 14 ; LR, ref 15 ; $A C$, ref $16 ; B A$, ref 17 ; AL, ref $18 ; B G$, ref $19 ; B O$, ref $20 ; L L$, ref 21 .
abilities in place of components $\mathbf{T}_{i j}$ when atom $i$ is bonded to $j$. In eq $10, \Delta \alpha_{i}$ is rotated along the direction of $i j$, and if $i$ is not bonded to $j$, its effect is attennated by a function of the distance between the atoms. The parameters $c=2.54$ and $d=0.25$ are found to refine the effect of the molecular environment on the anisotropic contributions to $\alpha_{i}$ and yield the present empirical formulation.

The point and modified dipole models ${ }^{1-9}$ were based on finding a set of atomic polarizabilities that yielded principal components of molecular polarizability in agreement with experimental results. The dipole approximation is valid when the separation between dipoles is large compared to the induced dipole moments. Thole ${ }^{9}$ analyzed this problem and noted that the polarizabilities, which are approximately the atomic volumes, could be used to design shape and size parameters in potential energy functions. He proposed several functional forms, which obey certain boundary conditions, and obtained good results.

In contrast, the present paper includes the effect of the dipole approximation with increasing distance, by omitting the terms with bonded atoms and assuming that these effects are replaced by atomic anisotropies. In the case of diatomic molecules, the anisotropies can be rationalized by using an atomic polarizability matrix that reproduces the molecular polarizability without the dipole tensor. These ideas are extended to the approximations for polyatomic molecules.

The neglect of nearest neighbor dipole terms on the polarizability tensor can be justified both qualitatively and quantitatively. Dipole expansions are valid for long-range interactions. Consequently, the adjustment of the atomic polarizabilities within the immediate molecular environment in the full dipole approximation ${ }^{7-9}$ must compensate for overemphasis of these short-range interactions. The present empirical method replaces these neighbor terms with local bond polarizability components, with the result that the atomic hybrid polarizabilities need no further adjustment. The rationale is based on the fact that atoms undergo hy bridization in the immediate bonding environment with virtually no influence from the nonbonded interactions. This strongly influences the polarizability parameters. For example, in an empirical analysis, the additivity of atomic hybrid components ${ }^{10.11}$ supports the neglect of nonbonded effects in the present model of average molecular polarizability, as demonstrated by the transferability of parameters among many molecules without regard to the atomic environment. ${ }^{10.11}$ The atom to which a given atom is bonded is not required, but only its state of hybridization. The same behavior appears valid for the components of molecular polarizability. The
present data set yields results within experimental error. Consequently, small modifications in the atomic hybrid polarizabilities will not improve the empirical fit consistently for the entire data set, because they depend on the number of data points used and the manner in which the data set is partitioned. Revisions made every time new data were added did not yield an overall improvement. Until experimental data accurate to $1 \%$ are available, further refinements in the atomic components will not yield meaningful results.

A quantum mechanical and quantitative evaluation, support, and justification is found in recent calculations by Garmer and Stevens ${ }^{22}$ on the transferability of molecular distributed polarizabilities. Their results on methane, ethane, propane, acetonitrile, methylamine, methanol, ethanol, and dimethyl ether support the present study. By using localized molecular orbital techniques, they partition the bond polarizabilities at the bond centroids for the purpose of developing a set of transferable bond polarizability tensors. They found that the bond centroids and the polarizability tensors for the C-H bonds are transferable with a standard deviation of 0.14 , which is of the order of $3-4 \%$ for the average polarizability. This is equivalent to the errors in the experimental measurements. Although an exact derivation or extensive quantum mechanical calculations are not available for a partitioning into either bond or atomic contributions, it appears that the results to date support the importance of the influence of the local environment on the polarizability with long-range dipolar adjustments to obtain the proper molecular polarizability components. The present empirical method yields results that are within experimental errors and internally consistent. Another important feature is that it is sufficiently simple to code and it is an effective approximation for molecular modeling.

Effects of nonadditivity begin to appear in large polymeric systems. The increase in the longitudinal polarizabilities of conjugated polymers is the basis for their utilization in optical signal processing devices. ${ }^{23}$ Theoretical longitudinal polarizabilities of polydiacetylene and polybutatriene have been calculated with the STO-3G and 4-31G basis sets by Kirtman and Hasan ${ }^{23}$ and with the INDO method by Kirtman. ${ }^{24}$ These calculations show that the increment by which the $z$ component changes reaches a plateau in the INDO but not in the STO basis sets. The increment, $\Delta \alpha_{z z} / N$ for the change in PDA, $\mathrm{C}_{6} \mathrm{H}_{6}+\mathrm{NC}_{2} \mathrm{H}_{2}$,

[^2]| molecule | method $^{\text {a }}$ | a. $\lambda^{3}$ | $f_{1}$ | $f_{2}$ | $f_{3}$ | $\begin{gathered} 1008 f_{6}^{b} \end{gathered}$ | molecule | method ${ }^{\text {a }}$ | $\boldsymbol{\alpha}, \AA^{3}$ | $f_{1}$ | $f_{2}$ | $f_{3}$ | $\begin{gathered} 100 \delta f_{0}^{b} \\ \% \end{gathered}$ | molecule | method $^{\text {a }}$ | $\boldsymbol{\alpha}, \AA^{3}$ | $f_{1}$ | $f_{2}$ | $f_{3}$ | $\begin{gathered} 1008 f^{b} \\ \hline \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| H2 | E | 0.79 | 0.304 | 0.304 | 0.392 |  | dimethyl ether | E | 5.24 | 0.279 | 0.314 | 0.406 |  | $\left(\mathrm{CH}_{3}\right)_{3} \mathrm{CCN}$ | E | 9.59 | 0.314 | 0.314 | 0.372 |  |
|  | A | 0.80 | 0.800 | 0.100 | 0.100 | 49.9 |  | A | 5.20 | 0.268 | 0.280 | 0.452 | 6.5 |  | A | 9.84 | 0.319 | 0.321 | 0.360 | 1.6 |
|  | B | 0.80 | 0.304 | 0.304 | 0.392 | 0.0 |  | B | 5.20 | 0.268 | 0.347 | 0.385 | 8.4 |  | B | 9.84 | 0.319 | 0.319 | 0.363 | 1.1 |
|  | T | 0.76 | 0.303 | 0.303 | 0.395 | 0.3 |  | T | 5.24 | 0.291 | 0.292 | 0.417 | 2.9 |  | T | 9.83 | 0.319 | 0.319 | 0.362 | 1.3 |
|  | M | 0.77 | 0.304 | 0.304 | 0.392 | 0.2 |  | M | 5.08 | 0.289 | 0.297 | 0.415 | 2.2 |  | M | 9.97 | 0.317 | 0.317 | 0.367 | 0.7 |
| N2 | E | 1.76 | 0.275 | 0.275 | 0.451 |  | $p$-dioxane | E | 8.60 | 0.271 | 0.364 | 0.364 |  | $\mathrm{CCl}_{3} \mathrm{CN}$ | E | 10.42 | 0.329 | 0.329 | 0.342 |  |
|  | A | 1.76 | 0.136 | 0.136 | 0.727 | 33.9 |  | A | 8.68 | 0.244 | 0.370 | 0.387 | 3.6 |  | A | 10.28 | 0.330 | 0.330 | 0.342 | 0.2 |
|  | B | 1.76 | 0.275 | 0.275 | 0.451 | 0.0 |  | B | 8.68 | 0.248 | 0.362 | 0.389 | 3.4 |  | M | 10.24 | 0.330 | 0.330 | 0.340 | 0.3 |
|  | T | 1.72 | 0.294 | 0.294 | 0.412 | 4.7 |  | T | 8.84 | 0.279 | 0.324 | 0.397 | 5.3 | $\mathrm{CH}_{3} \mathrm{~F}$ | E | 2.62 | 0.298 | 0.298 | 0.405 |  |
|  | M | 1.91 | 0.275 | 0.275 | 0.450 | 0.1 |  | M | 8.61 | 0.282 | 0.325 | 0.393 | 5.0 |  | A | 2.47 | 0.337 | 0.337 | 0.326 | 9.6 |
| 02 | E | 1.60 | 0.254 | 0.254 | 0.493 |  | methanol | E | 3.32 | 0.266 | 0.324 | 0.410 |  |  | M | 2.51 | 0.332 | 0.332 | 0.337 | 8.3 |
|  | A | 1.60 | 0.177 | 0.177 | 0.647 | 18.8 |  | A | 3.04 | 0.291 | 0.302 | 0.407 | 3.4 | $\mathrm{CHF}_{3}$ | E | 2.81 | 0.319 | 0.340 | 0.340 |  |
|  | B | 1.60 | 0.255 | 0.255 | 0.490 | 0.4 |  | B | 3.04 | 0.274 | 0.308 | 0.418 | 1.9 |  | A | 2.78 | 0.321 | 0.339 | 0.339 | 0.3 |
|  | T | 1.49 | 0.281 | 0.281 | 0.438 | 6.8 |  | T | 3.34 | 0.295 | 0.310 | 0.395 | 3.6 |  | M | 2.32 | 0.330 | 0.335 | 0.335 | 1.3 |
|  | M | 1.14 | 0.252 | 0.252 | 0.490 | 0.8 |  | M | 3.25 | 0.294 | 0.315 | 0.391 | 3.5 | $\mathrm{CH}_{3} \mathrm{Cl}$ | E | 4.55 | 0.292 | 0.292 | 0.416 |  |
| CO | E | 1.95 | 0.278 | 0.278 | 0.444 |  | ethanol | E | 5.08 | 0.295 | 0.327 | 0.378 |  |  | A | 4.54 | 0.267 | 0.267 | 0.467 | 1.7 |
|  | A | 1.95 | 0.138 | 0.138 | 0.723 | 34.2 |  | A | 5.11 | 0.256 | 0.280 | 0.463 | 10.4 |  | M | 4.54 | 0.285 | 0.285 | 0.431 | 1.7 |
|  | B | 1.95 | 0.277 | 0.277 | 0.445 | 0.2 |  | B | 5.11 | 0.286 | 0.321 | 0.393 | 1.8 | $\mathrm{CH}_{2} \mathrm{Cl}_{2}$ | E | 6.82 | 0.262 | 0.308 | 0.430 |  |
|  | T | 1.82 | 0.295 | 0.295 | 0.409 | 4.2 |  | T | 5.26 | 0.289 | 0.306 | 0.405 | 3.5 |  | A | 6.65 | 0.234 | 0.291 | 0.475 | 5.5 |
|  | M | 1.92 | 0.282 | 0.282 | 0.444 | 1.0 |  | M | 5.08 | 0.291 | 0.346 | 0.363 | 2.5 |  | M | 6.47 | 0.256 | 0.312 | 0.431 | 0.7 |
| ethane | E | 4.48 | 0.314 | 0.314 | 0.372 |  | formaldehyde | E | 2.45 | 0.249 | 0.376 | 0.376 |  | $\mathrm{CHCl}_{3}$ | E | 8.53 | 0.263 | 0.368 | 0.368 |  |
|  | A | 4.46 | 0.267 | 0.267 | 0.466 | 11.6 |  | A | 2.31 | 0.135 | 0.316 | 0.549 | 21.6 |  | A | 8.73 | 0.223 | 0.388 | 0.388 | 4.9 |
|  | B | 4.46 | 0.280 | 0.280 | 0.441 | 8.4 |  | B | 2.31 | 0.186 | 0.369 | 0.445 | 9.4 |  | M | 8.39 | 0.254 | 0.373 | 0.373 | 1.1 |
|  | $\mathrm{T}$ | 4.46 | 0.316 | 0.316 | 0.368 | 0.5 |  | T | 2.54 | 0.244 | 0.354 | 0.402 | 3.4 | $\mathrm{CH}_{3} \mathrm{Br}$ | E | 5.61 | 0.295 | 0.295 | 0.411 |  |
|  | M | 4.44 | 0.314 | 0.314 | 0.372 | 0.0 |  | M | 2.70 | 0.264 | 0.342 | 0.395 | 4.2 |  | A | 5.39 | 0.267 | 0.267 | 0.466 | 6.8 |
| propane | E | 6.38 | 0.300 | 0.300 | 0.400 |  | acetone | E | 6.39 | 0.254 | 0.372 | 0.373 |  |  | M | 5.24 | 0.294 | 0.294 | 0.412 | 0.2 |
|  | A | 6.58 | 0.257 | 0.289 | 0.453 | 6.8 |  | A | 6.46 | 0.263 | 0.368 | 0.370 |  | $\mathrm{CHBr}_{3}$ | E | 11.84 | 0.268 | 0.366 | 0.366 |  |
|  | B | 6.58 | 0.311 | 0.324 | 0.365 | 4.4 |  | B | 6.46 | 0.264 | 0.365 | 0.371 | 1.2 |  | A | 10.49 | 0.262 | 0.369 | 0.369 | 0.7 |
|  | T | 6.29 | 0.302 | 0.317 | 0.381 | 2.6 |  | T | 6.32 | 0.272 | 0.343 | 0.385 | 3.6 |  | M | 11.89 | 0.225 | 0.387 | 0.387 | 5.3 |
|  | M | 6.28 | 0.299 | 0.320 | 0.381 | 2.8 |  | M | 6.37 | 0.277 | 0.343 | 0.381 | 3.8 | $\mathrm{CH}_{3} \mathrm{I}$ | E | 7.59 | 0.302 | 0.302 | 0.396 |  |
| cyclopetane | E | 9.15 | 0.306 | 0.334 | 0.360 |  | acetonitrile |  | 4.48 | 0.286 | 0.286 | 0.427 |  |  | A | 7.29 | 0.269 | 0.269 | 0.463 | 8.2 |
|  | A | 9.05 | 0.264 | 0.367 | 0.369 | 5.5 |  | A | 4.14 | 0.232 | 0.232 | 0.536 | 13.4 |  | M | 7.64 | 0.305 | 0.305 | 0.391 | 0.7 |
|  | B | 9.05 | 0.277 | 0.358 | 0.365 | 3.8 |  | B | 4.14 | 0.263 | 0.263 | 0.474 | 5.7 | $\mathrm{CH}_{3} \mathrm{I}$ | E | 18.04 | 0.309 | 0.345 | 0.345 |  |
|  | T | 9.12 | 0.306 | 0.334 | 0.360 | 2.1 |  | T | 4.24 | 0.280 | 0.280 | 0.440 | 1.6 |  | A | 18.48 | 0.215 | 0.393 | 0.393 | 11.6 |
|  | M | 9.18 | 0.289 | 0.353 | 0.358 | 2.6 |  | M | 4.46 | 0.276 | 0.276 | 0.474 | 2.6 |  | M | 17.69 | 0.263 | 0.369 | 0.369 | 5.7 |
| cyclohexane | E | 11.00 | 0.284 | 0.358 | 0.358 |  |  |  |  |  |  |  |  |  |  |  |  |  |  |  |
|  | A | 10.95 | 0.254 | 0.373 | 0.373 | 3.7 |  |  |  |  |  |  |  |  |  |  |  |  |  |  |
|  | B | 10.95 | 0.267 | 0.367 | 0.367 | 2.1 |  |  |  |  |  |  |  |  |  |  |  |  |  |  |
|  | T | 10.95 | 0.288 | 0.288 | 0.356 | 5.0 |  |  |  |  |  |  |  |  |  |  |  |  |  |  |
|  | M | 11.01 | 0.284 | 0.358 | 0.358 | 0.0 |  |  |  |  |  |  |  |  |  |  |  |  |  |  |

${ }^{6} \delta f$ is defined by eq 20 .

Table 1V. Polarizability Components of Minimum-Energy Conformations of di-n-Propyl Ether

| ${ }_{\text {mol }}$ | $f_{1}$ | $f_{2}$ | $f_{3}$ | $\boldsymbol{\alpha}$ (ahp) | 1008f, \% | ref | U | conformation | $w$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 12.42 | 0.2667 | 0.2850 | 0.4482 | 13.73 | 10.3 | T | -6.4 | tott | 1 |
| 12.53 | 0.3113 | 0.3246 | 0.3645 |  |  | E |  |  |  |
| 12.42 | 0.2744 | 0.2995 | 0.4261 | 13.58 | 7.6 | T | -6.5 | $\mathbf{t u O t g}^{-}$ | 4 |
| 12.53 | 0.3113 | 0.3246 | 0.3645 |  |  | E |  |  |  |
| 12.42 | 0.2942 | 0.3064 | 0.3994 | 13.43 | 4.3 | T | -6.7 | gtotg ${ }^{-}$ | 2 |
| 12.53 | 0.3113 | 0.3246 | 0.3645 |  |  | E |  |  |  |
| 12.42 | 0.2760 | 0.3271 | 0.3969 | 13.41 | 4.8 | T | -6.8 | $\mathrm{gtOtg}{ }^{-}$ | 2 |
| 12.53 | 0.3113 | 0.3246 | 0.3645 |  |  | E |  |  |  |
| 12.4 | 0.2783 | 0.3055 | 0.4161 | 13.41 | 6.4 | T | -6.6 | average |  |

[^3]Table V. Molecular Polarizability with Successive Approximations to $\alpha_{\mathrm{A}}$ and $\mathrm{T}^{a}$

| molecule | $\alpha$ (ahp), $\AA^{3}$ | $f_{1}$ | $f_{2}$ | $f_{3}$ | $\alpha(\mathrm{mol}), \AA^{3}$ | 1008f, \% | approx |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| CS2 | 8.74 | 0.3333 | 0.3333 | 0.3333 | 8.74 | 29.9 | 1 |
|  | 8.74 | 0.2498 | 0.2498 | 0.5004 | 8.74 | 9.4 | 2 |
|  | 8.74 | 0.3045 | 0.3045 | 0.3910 | 8.88 | 22.8 | 3 |
|  | 8.74 | 0.2070 | 0.2070 | 0.5860 | 9.99 | 1.1 | 4 |
|  | 8.74 | 0.2113 | 0.2113 | 0.5774 |  |  | exp |
| ethane | 4.44 | 0.3333 | 0.3333 | 0.3333 | 4.44 | 9.2 | 1 |
|  | 4.44 | 0.3244 | 0.3244 | 0.3512 | 4.44 | 7.0 | 2 |
|  | 4.44 | 0.3108 | 0.3108 | 0.3784 | 4.55 | 3.7 | 3 |
|  | 4.44 | 0.3148 | 0.3148 | 0.3704 | 4.64 | 4.6 | 4 |
|  | 4.48 | 0.2961 | 0.2961 | 0.4085 |  |  | exp |
| nitrobenzene | 12.41 | 0.3333 | 0.3333 | 0.3333 | 12.41 | 18.3 | 1 |
|  | 12.41 | 0.2703 | 0.3633 | 0.3664 | 12.41 | 11.8 | 2 |
|  | 12.41 | 0.2275 | 0.3544 | 0.4181 | 13.45 | 5.0 | 3 |
|  | 12.41 | 0.1854 | 0.3769 | 0.4377 | 14.36 | 4.3 | 4 |
|  | 12.92 | 0.1999 | 0.3418 | 0.4582 |  |  | exp |
| diethyl ether | 8.75 | 0.3333 | 0.3333 | 0.3333 | 8.75 | 12.0 | 1 |
|  | 8.75 | 0.3280 | 0.3281 | 0.3439 | 8.75 | 10.7 | 2 |
|  | 8.75 | 0.2793 | 0.2996 | 0.4211 | 9.25 | 1.3 | 3 |
|  | 8.75 | 0.2783 | 0.2962 | 0.4254 | 9,50 | 1.0 | 4 |
|  | 8.73 | 0.2700 | 0.3005 | 0.4299 |  |  | exp |
| guanine | 14.59 | 0.3333 | 0.3333 | 0.3333 | 14.59 | 19.2 | 1 |
|  | 15.68 | 0.2832 | 0.3545 | 0.3624 | 15.68 | 13.1 | 2 |
|  | 14.59 | 0.2230 | 0.3436 | 0.4335 | 16.03 | 5.8 | 3 |
|  | 14.59 | 0.1755 | 0.3582 | 0.4663 | 17.39 | 3.7 | 4 |
|  | 13.60 | 0.1799 | 0.3819 | 0.4385 |  |  | $\exp$ |
| $\mathrm{CH}_{3} \mathrm{C}\left(\mathrm{CH}_{2} \mathrm{O}\right)_{3} \mathrm{PO}$ | 12.87 | 0.3333 | 0.3333 | 0.3333 | 12.87 | 23.2 | 1 |
|  | 12.87 | 0.3097 | 0.3097 | 0.3806 | 12.87 | 3.6 | 2 |
|  | 12.87 | 0.3076 | 0.3076 | 0.3846 | 13.54 | 0.5 | 3 |
|  | 12.87 | 0.2969 | 0.2969 | 0.4062 | 13.98 | 3.1 | 4 |
|  | 12.84 | 0.3097 | 0.3097 | 0.3806 |  |  | exp |
| acetonitrile | 4.46 | 0.3333 | 0.3333 | 0.3333 | 4.46 | 11.5 | 1 |
|  | 4.46 | 0.3122 | 0.3122 | 0.3756 | 4.53 | 6.3 | 2 |
|  | 4.46 | 0.2972 | 0.2972 | 0.4056 | 4.46 | 2.6 | 3 |
|  | 4.46 | 0.2769 | 0.2769 | 0.4462 | 4.73 | 2.4 | 4 |
|  | 4.48 | 0.2865 | 0.2865 | 0.4271 |  |  | exp |

${ }^{a}$ The four successive approximations are described in the text. $\alpha$ (ahp) is the average molecular or atomic hybrid polarizability. $f_{\mu}, \mu=1,2$, and 3 are the scaled principal components for each approximation. $\alpha(\mathrm{mol})$ is the nonscaled average molecular polarizability. $\delta f$ is defined by eq 20 .

Table VI. Atomic Polarizabilities for Various Methods ${ }^{\text {a }}$

| atom | atomic hybrid | dipole models A, B | additive model V | ahp model M | modified dipole model T |
| :---: | :---: | :---: | :---: | :---: | :---: |
| H(alkane) | H | 0.135 | 0.407 | 0.387 | 0.514 |
| H(alcohol) | H | 0.135 | 0.405 | 0.387 |  |
| H(aldehyde) | H | 0.167 |  | 0.387 |  |
| H(amide) | H | 0.161 |  | 0.387 |  |
| C(alkane) | CTE | 0.878 | 1.027 | 1.061 | 1.405 |
| C(Carbonyl) | CTR | 0.616 | 1.027 | 1.352 |  |
| C(nitrile) | CDI | 0.36 | 0.928 | 1.283 |  |
| $N$ (amide) | NPI2 | 0.530 |  | 1.090 | 1.105 |
| N(nitrile) | NDI | 0.52 | 1.236 | 0.956 |  |
| O(alcohol) | OTE | 0.465 | 0.604 | 0.637 | 0.862 |
| O (ether) | OTE | 0.465 | 0.651 | 0.637 |  |
| O(carbonyl) | OTR4 | 0.434 | 0.841 | 0.569 |  |
| F |  | 0.32 | 0.32 | 0.296 |  |
| Cl |  | 1.91 | 2.32 | 2.315 |  |
| Br |  | 2.88 | 3.465 | 3.013 |  |
| 1 |  | 4.69 | 5.531 | 5.415 |  |

increases by $21 \%$ for the INDO calculations and by $19 \%$ in sample calculations in the present unscaled modified dipole method.

Similar results are obtained for polybutatriene. The interpretation within the present modified dipole model is that the components $T_{\text {piqj }}$ contribute to the A matrix in eq 5 until the nearest, next nearest, etc. neighbor interactions become sufficiently damped so that the long-range contributions vanish. Then the $\mathbf{B}$ matrix in eqs 6 and 7 and finally the molecular polarizability matrix become saturated. The additivity and scaling methods used in the present semiempirical methods ${ }^{10,11}$ apply to small molecules and should be used accordingly. Sufficient data on linear polymeric systems are not available to provide additional calibrations for nonlinear effects, especially on the components of polarizability.

Essential features of the present method are as follows: Atomic anisotropies are introduced to ensure that the average atomic polarizabilities remained fixed at the $\alpha_{A}$ (ahp) values ${ }^{1}$ and the final molecular polarizability matrix is normalized and rescaled to yield $\alpha(\mathrm{ahp})$ and the principal components of molecular polarizability. The present approach unifies the additivity of atomic polarizabilities of the ahp method with a modified dipole tensor approximation.

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[^3]:    of is defined by eq 20

